X. Investigation of an Extensive Class of Partial Differential Equations of the Second Order, in which the Equation of Laplace's Functions is included.

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Theorem. If u be a function of x and y satisfying the equation

$$\frac{d^2u}{dxdy} + \alpha_n e^{\varphi} u = 0,$$

where

$$\frac{d^2\varphi}{dxdy} + ce^{\varphi} = 0,$$

then the solution will be

$$u = D^{-n}v_n$$

where

$$\mathbf{D} = e^{-\varphi} \frac{d}{dy},$$

where

$$v_n = \int e^{-\frac{\beta_n}{\Delta\beta_n}} \varphi_{\chi} y dy + \psi x,$$

 χy and ψx arbitrary functions of y and x,

and

$$\beta_n = \frac{\Delta \alpha_{n-1} \cdot \Delta \alpha_{n-2} \cdot \ldots}{(\Delta \alpha_{n-1} - c)(\Delta \alpha_{n-2} - c) \cdot \ldots},$$

where

$$\Delta \alpha_r = \alpha_{r+1} - \alpha_r$$

and α_n is a function of *n* vanishing for n=0 and for n=-1.

I will proceed to demonstrate this curious theorem as briefly as possible.

According to the notation, we may write the given equation

$$D\frac{du}{dx} + \alpha_n u = 0.$$

Then if

$$v_n = \mathbf{D}^n u$$

we have

$$v_{n+1} = \mathbf{D}^{n+1} u = \mathbf{D} v_n.$$

Suppose that

$$D\frac{dv_n}{dx} + \beta_n z Dv_n = D^n \left\{ D\frac{du}{dx} + \alpha_n u \right\}. \qquad (\alpha.)$$

where z is a function of x and y to be determined, also β_n a function of u.

Writing n+1 for n in equation (α) , we ought to have

$$D\frac{dv_{n+1}}{dx} + \beta_{n+1}zDv_{n+1} = D^{n+1}\left\{D\frac{du}{dx} + \alpha_{n+1}u\right\} (\beta.$$

This circumstance will serve to determine z and β_n as follows: we have identically

$$\mathbf{D}^{n+1} \Big\{ \mathbf{D} \frac{du}{dx} + \alpha_{n+1} u \Big\} = \mathbf{D}^{n+1} \Big\{ \mathbf{D} \frac{du}{dx} + \alpha_{n} u \Big\} + \Delta \alpha_{n} \mathbf{D}^{n+1} u \Big\}$$

$$= \mathbf{D}^{2} \frac{dv_{u}}{dx} + \beta_{n} \mathbf{D} \Big\{ z \mathbf{D} v_{n} \Big\} + \Delta \alpha_{n} \cdot \mathbf{D} v_{n} \text{ by } (\alpha)$$

$$= \mathbf{D}^{2} \frac{dv_{n}}{dx} + \beta_{n} \mathbf{D} \Big\{ z v_{n+1} \Big\} + \Delta \alpha_{n} \cdot v_{n+1}.$$
But
$$\frac{dv_{n+1}}{dx} = \frac{d}{dx} \mathbf{D} v_{n} = \frac{d}{dx} \Big(e^{-\varphi} \frac{dv_{n}}{dy} \Big)$$

$$= -\frac{d\varphi}{dx} e^{-\varphi} \frac{dv_{n}}{dy} + e^{-\varphi} \frac{d}{dy} \cdot \frac{dv_{n}}{dx}$$

$$= -\frac{d\varphi}{dx} \mathbf{D} v_{n} + \mathbf{D} \frac{dv_{n}}{dx};$$

$$\therefore \mathbf{D}^{2} \frac{dv_{n}}{dx} = \mathbf{D} \Big\{ \frac{dv_{n+1}}{dx} + \frac{d\varphi}{dx} v_{n+1} \Big\}.$$
Hence
$$\mathbf{D} \frac{dv_{n+1}}{dx} + \mathbf{D} \Big(\frac{d\varphi}{dx} v_{n+1} \Big) + \beta_{n} \mathbf{D} (z v_{n+1}) + \Delta \alpha_{n} v_{n+1}.$$

ought to be identical with $\mathbf{D} \frac{dv_{n+1}}{dx} + \beta_{n+1} z \mathbf{D} v_{n+1}$, and hence the conditions

$$\frac{d\phi}{dx} + \beta_n z = \beta_{n+1} z$$

$$\mathbf{D} \left\{ \frac{d\phi}{dx} + \beta_n z \right\} = -\Delta \alpha_n.$$

Eliminating z, we have

or
$$\begin{aligned} \mathbf{D} \frac{d\varphi}{dx} &= -\frac{\Delta \alpha_n \Delta \beta_n}{\beta_{n+1}} = -c, \\ \frac{d^2 \varphi}{dx dy} + c e^{\varphi} &= 0, \\ \text{and} \qquad c \beta_{n+1} &= \Delta \alpha_n \{ \beta_{n+1} - \beta_n \}, \\ \text{or} \qquad \beta_{n+1} &= \frac{\Delta \alpha_n}{\Delta \alpha_n - c} \cdot \beta_n; \\ &\therefore \qquad \beta_n &= \frac{\Delta \alpha_{n-1} \cdot \Delta \alpha_{n-2} \cdot \dots}{(\Delta \alpha_{n-1} - c)(\Delta \alpha_{n-2} - c) \cdot \dots} \end{aligned}$$

by these determinations we establish the formula (β) as a consequence of (α) , and therefore if the formula (α) be true for any value of n, it will be (subject to the above conditions) true for the next superior value.

Now, when n=0, $v_0=D^0u=u$, and provided α_0 and α_{-1} are each =0, α_0 and $\Delta\alpha_{-1}$ will be each 0, and α_0 and β_0 each =0, and the equation (α .) reduces to $D \frac{dv_0}{dx}$ =D $\frac{du}{dx}$, and is therefore true for u=0. Under these restrictions it will therefore be true for any positive integral value of n. Now the symbol D represents $e^{-\phi} \frac{d}{dy}$, and therefore if U=0, DⁿU=0, so that we have

or
$$\begin{aligned} \mathbf{D} \frac{dv_n}{dx} + \beta_n z \mathbf{D} v_n &= 0, \\ e^{-\varphi} \frac{d^2 v_n}{dx dy} + \frac{\beta_n}{\Delta \beta_n} \cdot \frac{d\varphi}{dx} \cdot e^{-\varphi} \frac{dv_n}{dy} &= 0, \\ \\ \frac{\frac{d}{dx} \cdot \frac{dv_n}{dy}}{\frac{dv_n}{dy}} &= -\frac{\beta_n}{\Delta \beta_n} \cdot \frac{d\varphi}{dx}; \end{aligned}$$

 \therefore integrating with respect to x,

$$\frac{dv_n}{dy} = e^{-\frac{\beta_n}{\Delta\beta_n}} \varphi \chi y$$

$$v_n = \int e^{-\frac{\beta_n}{\Delta\beta_n}} \varphi \chi y dy + \psi x,$$

$$u = \mathbf{D}^{-n} v_n = \int e^{\varphi} \int e^{\varphi} \dots v_n dy dy \dots$$

and

the integral sign repeated n times. The theorem is therefore demonstrated.

It may be easily shown that the equation of LAPLACE's coefficients is included in the class here considered.

The equation of Laplace by a proper choice of independent variables assumes the form

$$\frac{d^2u}{dxdy} + \frac{n \cdot n + 1}{4\cos^2 \frac{y - x}{2}} \cdot u = 0.$$

Hence with reference to the preceding investigation,

$$D = \cos^2 \frac{y - x}{2} \cdot \frac{d}{dy} \text{ and } \alpha_n = \frac{n \cdot \overline{n+1}}{4}.$$
Hence
$$e^{-\varphi} = \cos^2 \frac{y - x}{2};$$

$$\therefore \qquad \frac{d\varphi}{dx} = -\tan \frac{y - x}{2}$$

$$\frac{d^2\varphi}{dxdy} + \frac{1}{2}e^{\varphi} = 0.$$

Hence $c = \frac{1}{2}$ Also $\Delta \alpha_n = \frac{n+1}{2}$; $\Delta \alpha_n - c = \frac{n}{2}$;

$$\beta_n = n \text{ and } \Delta \beta_n = 1.$$

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Inserting these values in the final formula, we have

$$v_n = \int \cos^{2n} \frac{y-x}{2} \chi y dy + \psi x,$$

and

$$u = \int \cos^{-2} \frac{y - x}{2} \int \cos^{-2} \frac{y - x}{2} \dots v_n dy dy \dots n \text{ times,}$$

which agrees with Mr. Hargreave's solution.